

An Estimator for statistical anisotropy from the CMB bispectrum

(based on N.Bartolo, E.D, M.Liguori S.Matarrese and A.Riotto,
to appear)

Emanuela Dimastrogiovanni

Ann Arbor, 05/14/2011

Statistical anisotropy

- Observational hints of statistical anisotropy from Cosmic Microwave Background (CMB) data analysis
(e.g. cold spots; low amplitude of the quadrupole moment; alignment between quadrupole and octupole; lack of large angular scale CMB power; dipolar and quadrupolar power asymmetry.)
- Primordial vector field models of inflation predicting statistical anisotropy in the power spectrum and in higher order cosmological correlators
(e.g. hybrid inflation, vector curvaton, $f(\phi)$ models.)

$$P(\vec{k}) = P^{iso}(k) \left[1 + g(k) (\hat{k} \cdot \hat{n})^2 \right] \rightarrow g = 0.29 \pm 0.031, \parallel \text{ecliptic poles}$$

(astro-ph/0701357, arXiv:0709.1144) (arXiv : 0807.2242, 0908.0963, 0911.0150)

Complementary bispectrum analysis:

$$\begin{aligned} B(\vec{k}_1, \vec{k}_2, \vec{k}_3) = & B^{iso}(k_1, k_2) \left[1 + \Gamma (\hat{k}_1 \cdot \hat{N})^2 + \Delta (\hat{k}_2 \cdot \hat{N})^2 + \Theta (\hat{k}_1 \cdot \hat{N})^2 (\hat{k}_2 \cdot \hat{N})^2 \right. \\ & \left. + \Omega (\hat{k}_1 \cdot \hat{k}_2) (\hat{k}_1 \cdot \hat{N}) (\hat{k}_2 \cdot \hat{N}) \right] + 2 \text{ perms.} \end{aligned}$$

It turns out that: there are some models where the anisotropy of the power spectrum is not observable, whereas the anisotropic amplitude of the bispectrum can be relatively large!

Anisotropic CMB bispectrum

$$B_{m_1 m_2 m_3}^{l_1 l_2 l_3} \equiv \langle a_{l_1 m_1} a_{l_2 m_2} a_{l_3 m_3} \rangle = (4\pi)^3 (-i)^{l_1 + l_2 + l_3} \int \frac{d^3 k_1}{(2\pi)^3} \frac{d^3 k_2}{(2\pi)^3} \frac{d^3 k_3}{(2\pi)^3} \Delta_{l_1}(k_1) \Delta_{l_2}(k_2) \Delta_{l_3}(k_3) \\ \times Y_{l_1 m_1}(\hat{k}_1) Y_{l_2 m_2}(\hat{k}_2) Y_{l_3 m_3}(\hat{k}_3) \langle \Phi(\vec{k}_1) \Phi(\vec{k}_2) \Phi(\vec{k}_3) \rangle$$

- Spherical harmonics expansion of temperature anisotropies:

$$\frac{\Delta T(\hat{n})}{T} = \sum_{lm} a_{lm} Y_{lm}(\hat{n})$$

$$a_{lm} = 4\pi (-i)^l \int \frac{d^3 k}{(2\pi)^3} \Delta_l(k) \Phi(\vec{k}) Y_{lm}(\hat{k})$$

- Primordial bispectrum:

$$\langle \zeta(\vec{k}_1) \zeta(\vec{k}_2) \zeta(\vec{k}_3) \rangle = (2\pi)^3 \delta^{(3)}(\vec{k}_1 + \vec{k}_2 + \vec{k}_3) B(\vec{k}_1, \vec{k}_2, \vec{k}_3) = \left(\frac{5}{3}\right)^3 \langle \Phi(\vec{k}_1) \Phi(\vec{k}_2) \Phi(\vec{k}_3) \rangle$$

- Simplest template:

$$B(\vec{k}_1, \vec{k}_2, \vec{k}_3) = B^{iso}(k_1, k_2) \left[1 + g_1 (\hat{k}_1 \cdot \hat{n})^2 + g_2 (\hat{k}_2 \cdot \hat{n})^2 \right] + perms.$$

$$\rightarrow g_1 (\hat{k}_1 \cdot \hat{n})^2 = \sum_{LM} \lambda_{LM} Y_{LM}(\hat{k}_1)$$

⇓

$$B_{m_1 m_2 m_3}^{l_1 l_2 l_3} = B_{m_1 m_2 m_3}^{l_1 l_2 l_3(I)} + B_{m_1 m_2 m_3}^{l_1 l_2 l_3(A)} = f_{NL} \left(B_{m_1 m_2 m_3}^{l_1 l_2 l_3(I)}|_{f_{NL}=1} + \sum_{L \geq 2, M} \lambda_{LM} B_{m_1 m_2 m_3}^{l_1 l_2 l_3(A)LM}|_{f_{NL}=1} \right)$$

from isotropic and **anisotropic** contributions:

$$B_{m_1 m_2 m_3}^{l_1 l_2 l_3(I)} = b_{l_1 l_2 l_3} G_{m_1 m_2 m_3}^{l_1 l_2 l_3} \equiv \left(\frac{3}{5}\right)^3 \left(\frac{2}{\pi}\right)^3 \int dx x^2 \int dk_1 dk_2 dk_3 (k_1 k_2 k_3)^2 \Delta_{l_1}(k_1) \Delta_{l_2}(k_2) \Delta_{l_3}(k_3) \\ \times j_{l_1}(k_1 x) j_{l_2}(k_2 x) j_{l_3}(k_3 x) B^{iso}(k_1, k_2, k_3) G_{m_1 m_2 m_3}^{l_1 l_2 l_3}$$

$$B_{m_1 m_2 m_3}^{l_1 l_2 l_3(A)LM} = \left(\frac{3 \times 6}{5\pi}\right)^3 \sum_{l'_1 m'_1} G_{m'_1 m_2 m_3}^{l'_1 l_2 l_3} G_{m_1 - m'_1 M}^{l_1 l'_1 L} (-1)^{l'_1} (i)^{l_1 + l'_1} (-1)^{m'_1} b_{l'_1 l_2 l_3}^{l_1 l'_1} + 2 perms.$$

Results and conclusions

Estimators and Fisher matrix for our anisotropy parameters:

$$\hat{\lambda}_{LM} = \frac{1}{F_{\lambda_{LM}\lambda_{LM}}} \left(f_{NL} B^{(A)LM} - f_{NL}^2 B^{(I)} B^{(A)LM} \right)$$

$$F_{\lambda_{LM}\lambda_{LM}} = f_{NL}^2 \left\langle B^{(A)LM} B^{(A)L-M} (-1)^M \right\rangle$$

$$B^{(I)} \equiv \frac{1}{6} \sum_{l_1 m_j} B_{m_1 m_2 m_3}^{l_1 l_2 l_3 (I)} |_{f_{NL}=1} \left(\frac{a_{l_1 m_1}^* a_{l_2 m_2}^* a_{l_3 m_3}^*}{c_{l_1} c_{l_2} c_{l_3}} - \frac{(-1)^{m_2}}{c_{l_1} c_{l_2}} \delta_{l_2 l_3} \delta_{m_2 - m_3} a_{l_1 m_1}^* - 2 \text{ perms.} \right)$$

$$B^{(A)LM} \equiv \frac{1}{6} \sum_{l_1 m_j} B_{m_1 m_2 m_3}^{l_1 l_2 l_3 (A)LM} |_{f_{NL}=1} \left(\frac{a_{l_1 m_1}^* a_{l_2 m_2}^* a_{l_3 m_3}^*}{c_{l_1} c_{l_2} c_{l_3}} - \frac{(-1)^{m_2}}{c_{l_1} c_{l_2}} \delta_{l_2 l_3} \delta_{m_2 - m_3} a_{l_1 m_1}^* - 2 \text{ perms.} \right)$$

1 σ error for the quadrupole statistical anisotropy in the bispectrum (preliminary results):

- $\frac{1}{\sigma_{\lambda_{2M}}} \simeq 0.1 \frac{l^5 b_{III}}{(l^2 C_l)^{3/2}} \simeq 0.77 f_{NL} \left(\frac{l}{2000} \right)$
- $\sigma_{\lambda_{2M}} \simeq 0.04$ for $f_{NL} = 32$ and $l_{max} \simeq 2000$.

→ Primordial vector field models exist that predict a bispectrum statistical anisotropy that might be large enough to be within the sensitivity ranges of forthcoming experiments. However, a careful analysis of systematic errors and of other possible cosmological signals that can mimic statistical anisotropy will also be needed!